Rational solution of Painleve IV

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Asymptotic and Computational Aspects of Complex ODEs CRM PISA



FCT Fundação para a Ciência e a Tecnologia

MINISTÉRIO DA CIÊNCIA, TECNOLOGIA E ENSINO SUPERIOR

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Talk based on:

- A. Eremenko, D. Masoero *Poles of rational solutions of Painleve II* (unpublished 2013)
- D. Masoero, P. Roffelsen : *Poles and Zeroes of rational solutions of Painleve IV* (work in progress)



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- Painleve equations can be realised as isomonodromic deformation of a linear system Ψ' = A(y(z), y'(z), z; θ)Ψ

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- Conversely, if f is any meromorphic function with a a finite number of singularities then f = ψ₁/ψ₂ where the ψ_{1,2} solve a linear ODE ψ_{1,2}" = Pψ_{1,2} for some P as above [Nevanlinna, Elfving]

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- Conversely, if f is any meromorphic function with a a finite number of singularities then $f = \frac{\psi_1}{\psi_2}$ where the $\psi_{1,2}$ solve a linear ODE $\psi_{1,2}'' = P\psi_{1,2}$ for some P as above [Nevanlinna, Elfving]
- Topological classification: two meromorphic functions are equivalent iff they are topologically equivalent as covering maps of the sphere.

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- The 'values' of the transcendental singularities of *f* can be read off the Stokes multipliers [M.]
- The location of the (movable) critical points of *f* are related to the values of *y*.
- Poles of y actually simplifies the monodromy problem!
- Consequences: A-priori characterization of the set y⁻¹(w) by combinatorics (topology) → simple proof of surjectivity of the monodromy map (ex: PI).

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Some background. PIV

• Main example for this talk: Painleve IV

$$y_{zz}=\frac{1}{2y}y_z^2+\frac{3}{2}y^3+4zy^2+2(z^2+1-2\theta_\infty)y-\frac{8\theta^2}{y},\quad \theta,\theta_\infty\in\mathbb{C}.$$

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- There are at least three different isomonodromic representations for PIV. I use here the Garnier-Jimbo-Miwa: one fuchsian singularity at the origin and one irregular singularity at ∞.
- The monodromy map is rather involved. The resonant case is considered in [Kapaev, unpublished].

Rational solutions of PIV

 P_{IV} has a rational solution iff either

Hermite
$$heta = \frac{m}{2} + n, heta = \frac{m}{2},$$

or

Okamoto
$$heta=rac{m}{2}+n\pmrac{2}{3}, heta=rac{m}{2}$$

for some $m, n \in \mathbb{Z}$, furthermore for any such parameter values the associated rational solution is unique.

Hermite solutions

Three types of Hermite solutions

$$\omega_{m,n}^{(I)} = \frac{H'_{m+1,n}}{H_{m+1,n}} - \frac{H'_{m,n}}{H_{m,n}},$$

$$\omega_{m,n}^{(II)} = \frac{H'_{m,n}}{H_{m,n}} - \frac{H'_{m,n+1}}{H_{m,n+1}},$$

$$\omega_{m,n}^{(III)} = -2z + \frac{H'_{m,n+1}}{H_{m,n+1}} - \frac{H'_{m+1,n}}{H_{m+1,n}},$$

 $H_{m,n}(z)$ is the generalized Hermite polynomial.

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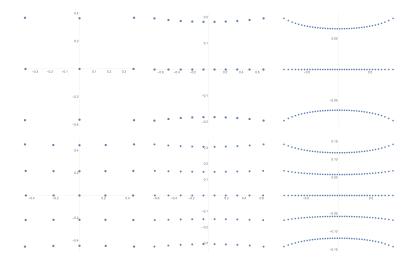
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 $H_{m,n}(z)$ is the generalized Hermite polynomial. Note: poles (residue ± 1) and zeroes coincide with zeroes of Hermite polynomials.

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Zeroes of generalized Hermite polynomials $H_{(m,n)}$



Problem [P. Clarkson]: explain the pictures!

D. Masoero (Universidade de Lisboa)

Pisa, CRM 9 / 26

Okamoto solutions

For $m, n \in \mathbb{Z}$,

$$\begin{split} \widetilde{\omega}_{m,n}^{(I)} &= -\frac{2}{3}z + \frac{Q'_{m+1,n}}{Q_{m+1,n}} - \frac{Q'_{m,n}}{Q_{m,n}}, \quad (1a) \\ \widetilde{\omega}_{m,n}^{(II)} &= -\frac{2}{3}z + \frac{Q'_{m,n}}{Q_{m,n}} - \frac{Q'_{m,n+1}}{Q_{m,n+1}}, \quad (1b) \\ \widetilde{\omega}_{m,n}^{(III)} &= -\frac{2}{3}z + \frac{Q'_{m,n+1}}{Q_{m,n+1}} - \frac{Q'_{m+1,n}}{Q_{m+1,n}}, \quad (1c) \end{split}$$

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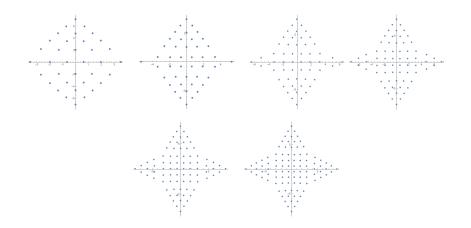
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Note: poles have residue ± 1 and coincide with zeroes of Okamoto polynomials.

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Rational Solutions of Painleve IV

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- Practice: Asymptotic analysis of these problems

Remark: 3 families of Hermite solutions \rightarrow 3 'inequivalent' monodromy problems for the same zeroes of the Hermite polynomials.

Isomonodromic equation at a pole. Hermite-I case.

Suppose $a \in \mathbb{C}$ is a pole of residue -1 of $\omega_{m,n}^{(I)}$ and

$$\omega_{m,n}^{(I)}(z) = \frac{-1}{z-a} + \cdots + b(z-a)^2 + \ldots$$

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The isomonodromic equation evaluated at z = a is

$$\psi''(\lambda) = \left((\lambda + a)^2 - (2m + n) - \frac{\kappa(a, b)}{\lambda} + \frac{n^2 - 1}{\lambda^2}\right)\psi(\lambda)$$

- The resonant singularity at $\lambda = 0$ is apparent: no logarithms.
- The Stokes multiplier σ_1 vanishes: $\psi_0 \sim \psi_2$.
- The subdominant solution ψ_0 vanishes at zero and has m-1 further simple zeroes.

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- ② f has 4 transcendental singularities but 2 of the 4 asymptotic values coincide: w₀, w₁, w₋₁, w₂ with w₂ = w₀ (↔ σ₁ = 0).

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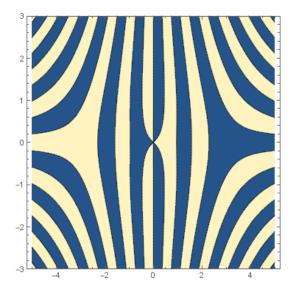
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A pictorial example. f for m = n = 3



The point $a \in \mathbb{C}$ is a pole of residue -1 of $\omega_{m,n}^{(I)}$

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- **2** The solution ψ_0 vanishes at $\lambda = 0$.

Indeed the further conditions

- $\psi(0)$ has m-1 simple zeroes.
- The Stokes multiplier σ_1 vanishes: $\psi_0 \sim \psi_2$

are satisfied because of the parametrization of the potential by m, n.

Asymptotic analysis

On the equation

$$\psi''(\lambda) = \left((\lambda + a)^2 - (2m + n) - \frac{\kappa(a, b)}{\lambda} + \frac{n^2 - 1}{\lambda^2}\right)\psi(\lambda)$$

we must impose two conditions in order to find the poles

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First step: For any fixed α , we look for κ such that the fuchsian singularity of

$$\psi''(\lambda) = \left(\lambda^2 + 2\alpha E^{\frac{1}{2}} - (1 - \alpha^2)E - \frac{\kappa}{\lambda} + \frac{n^2 - 1}{\lambda^2}\right)\psi(\lambda)$$

is apparent. A complicated algebraic equation of order n for κ .

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We compute the location of zeroes of Hermite polynomials in the limit $E = 2m + n \rightarrow \infty$ with *n* bounded.

We define $\alpha = E^{-\frac{1}{2}}a$ and restrict to $|1 - \alpha^2| > c > 0$.

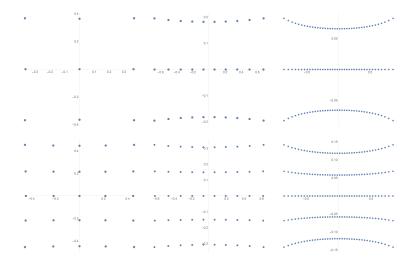
First step: For any fixed α , we look for κ such that the fuchsian singularity of

$$\psi''(\lambda) = \left(\lambda^2 + 2\alpha E^{\frac{1}{2}} - (1 - \alpha^2)E - \frac{\kappa}{\lambda} + \frac{n^2 - 1}{\lambda^2}\right)\psi(\lambda)$$

is apparent. A complicated algebraic equation of order *n* for κ . For $E \to +\infty$, we find:

$$\kappa=\sqrt{-1}j(1-lpha^2)^{rac{1}{2}}E^{rac{1}{2}}+O(E^{-rac{1}{2}})$$
 with $|j|\leq rac{n-1}{2}$ integer .

Zeroes of generalized Hermite polynomials $H_{(m,n)}$



In the pictures n = 3, 5. The integer *j* parametrizes the lines of poles!

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WKB expansion of ψ_0

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We have to find α such that the solution ψ_0 vanish at 0 For $O(E^{-\frac{1}{2}+\varepsilon}) \lesssim \lambda \lesssim (1-\alpha^2)E^{\frac{1}{2}}$, the WKB asymptotic holds

$$\begin{split} \psi_0 &\sim \sin\{\frac{\pi}{4} - (1 - \alpha^2)^{\frac{1}{2}} E^{\frac{1}{2}} \lambda - \frac{\sqrt{-1}j}{2} \log \lambda + \\ &+ \frac{E}{2} \left(-\alpha \sqrt{1 - \alpha^2} + \arccos \alpha \right) + \sqrt{-1} \frac{j}{2} \log 2(1 - \alpha^2) E^{\frac{1}{2}}) \} \end{split}$$

Asymptotic analysis of ψ_+

For $\lambda \rightarrow 0$ one can transform the original equation into a Whittaker equation

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Asymptotic analysis of ψ_+

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The subdominant solution ψ_+ vanishing at 0 has the asymptotic

$$\psi_+(\lambda) \sim M\left[rac{j}{2}, rac{n}{2}, 2\sqrt{-1}(1-lpha^2)^{rac{1}{2}} E^{rac{1}{2}}\lambda
ight], \ \lambda = O(E^{-\gamma}), \ \gamma > 0$$

where $M\left[\frac{j}{2},\frac{n}{2},\mu\right]$ is the Whittaker function.

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Solving the approximate equation, we find the α up to an error $O(E^{-\frac{4}{3}})$. Since the spacing among them is $O(E^{-1})$, all roots such that $|1 - \alpha^2| > c > 0$ are resolved.

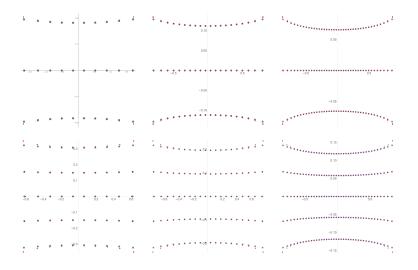
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We also computed the behavior for $\alpha \rightarrow \pm 1$ but there is no simple formula.

Asymptotics Vs. Numerics.



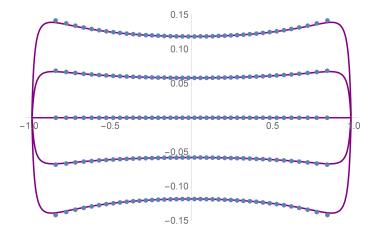
D. Masoero (Universidade de Lisboa)

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Layers of zero

In the large *m* limit, zeroes condensate on the curves $\Im \alpha = f_{i,n,m}(\Re \alpha)$



- We charatcterised poles and zero of Hermite solutions and computed its asymptotic distribution for $m \to \infty$ and n bounded (in case $|1 \alpha^2| > 0$ and $1 \alpha^2 \sim 0$).
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- Rational solutions and finite families of meromorphic functions: Is there any reason for this phenomenon?

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MANY THANKS FOR THE ATTENTION!

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